

Extended TQFT by Induced Representations

Jeffrey C. Morton

Instituto Tecnico Superior, Lisbon

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Definition

A Topological Quantum Field Theory is a monoidal functor:

$$Z : \mathbf{nCob} \rightarrow \mathbf{Vect}$$

An Extended TQFT (ETQFT) is a (weak) monoidal 2-functor

$$Z : \mathbf{nCob}_2 \rightarrow \mathbf{2Vect}$$

where \mathbf{nCob}_2 has

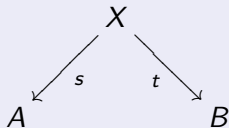
- **Objects:** $(n - 2)$ -dimensional manifolds
- **Morphisms:** $(n - 1)$ -dimensional cobordisms (manifolds with boundary, with ∂M a union of source and target objects)
- **2-Morphisms:** n -dimensional cobordisms with corners

We'll construct an ETQFT by factoring through a 2-category $\mathit{Span}(\mathbf{Gpd})$, then applying some universal process.

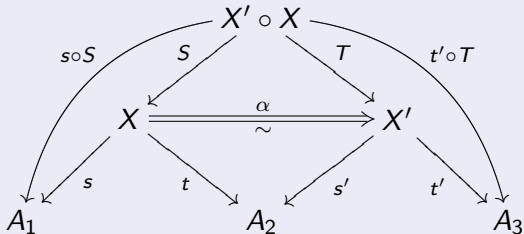
Definition (Part 1)

The bicategory $\text{Span}_2(\mathbf{Gpd})$ has:

- **Objects:** Groupoids
- **Morphisms:** Spans of groupoids:

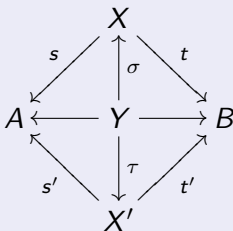


- Composition defined by *weak* pullback:



Definition (Part 2)

The **2-morphisms** of $\text{Span}_2(\mathbf{Gpd})$ are spans of *span maps*, commuting up to 2-cells of \mathbf{Gpd} :



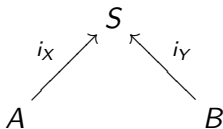
Composition is by weak pullback taken up to isomorphism.

Theorem

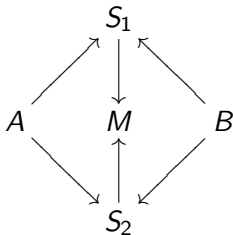
There is a monoidal structure on $\text{Span}_2(\mathbf{Gpd})$ induced by the product in \mathbf{Gpd} , with monoidal unit 1.

(Note: Roughly, $\text{Span}_2(\mathbf{C})$ will be the universal 2-category containing \mathbf{C} in which morphisms have ambidextrous adjoints.)

Cobordisms can be seen as *cospan*s of manifolds, with inclusions:

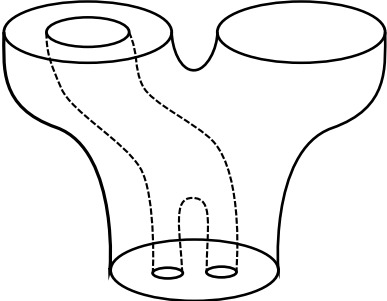


A cobordism between two cobordisms is a cospan of cospan maps:



(Note there are complications due to the fact that \mathbf{nCob}_2 is a *cubical* weak 2-category.)

Cobordisms in \mathbf{nCob}_2 actually live in $\text{Span}^2(\mathbf{ManCorn})$, as double cospans (here $n = 3$):

| \mathbf{nCob}_2 | $\text{Span}^2(\mathbf{ManCorn})$ |
|----------------------------------------------------------------------------------|---------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------|
|  | $ \begin{array}{ccccc} S^1 & \xrightarrow{i_A} & (A \amalg D) & \xleftarrow{i'_A \otimes i_D} & S^1 \amalg S^1 \\ i_1 \downarrow & & \downarrow \iota_1 & & \downarrow i_2 \\ Y & \xrightarrow{\iota_3} & M & \xleftarrow{\iota_4} & Y \\ i_2 \uparrow & & \uparrow \iota_2 & & \uparrow i_1 \\ S^1 \amalg S^1 & \xrightarrow{i_2} & Y & \xleftarrow{i_1} & S^1 \end{array} $ |

These form a “double bicategory”, which **induces** a bicategory since horizontal and vertical morphisms are composable.

For finite gauge group G , we get a functor:

$$\mathcal{A}_G : \mathbf{nCob}_2 \rightarrow \mathit{Span}(\mathbf{Gpd})$$

A *flat G -connection* on a manifold M is an object in

$$\mathcal{A}_G(M) = \mathit{hom}(\Pi_1(M), G)$$

Gauge transformations are natural transformations between these (giving a group element at each point).

Definition

Moduli space for *gauge theory*, for (finite) gauge group G . Given M , the groupoid $\mathcal{A}_G(M) = \mathit{Fun}(\pi_1(M), G)$ has:

- **Objects:** Flat connections on M (functors)
- **Morphisms** Gauge transformations (natural transformations)

(“Secretly” the groupoid is representing a *stack*. This is a standard situation for moduli spaces supporting symmetries.)

A connection on the cobordism $S : X \rightarrow Y$ in \mathbf{nCob}_2 can be pulled back along boundary inclusions by $(i_X)^*$ and $(i_Y)^*$, hence there is a span of the groupoids of flat connections:

$$\begin{array}{ccc} & \mathcal{A}_G(S) & \\ (i_X)^* \swarrow & & \searrow (i_Y)^* \\ \mathcal{A}_G(X) & & \mathcal{A}_G(Y) \end{array}$$

Theorem

$\mathcal{A}_G(-)$ defines a contravariant functor $\mathbf{ManCorn} \rightarrow \mathbf{Gpd}$, and a covariant functor $\mathbf{nCob}_2 \rightarrow \mathbf{Span}(\mathbf{Gpd})$.

Think of $\mathcal{A}_G(S)$ as a space (*stack*) of *histories*; intuitively s and t pick the starting and terminating *configuration* in A and B - compatible with gauge symmetry.

Goal: Using the induced 2-functor $\mathcal{A}_G(-) : \mathbf{nCob}_2 \rightarrow \mathit{Span}_2(\mathbf{Gpd})$, get an ETQFT $Z_G : \mathbf{nCob}_2 \rightarrow \mathbf{2Vect}$.

Theorem

There is a 2-functor (“2-linearization”):

$$\Lambda : \mathit{Span}_2(\mathbf{Gpd}) \rightarrow \mathbf{2Vect}$$

Where, recall:

Definition

$\mathbf{2Vect}$ is the 2-category of 2-vector spaces, which consists of:

- Objects: \mathbb{C} -linear abelian category, generated by simple objects
- Morphisms: **2-linear maps**: \mathbb{C} -linear (hence abelian) functor.
- 2-Morphisms: Natural transformations

Finite dimensional 2-vector spaces all look like \mathbf{Vect}^k , and 2-linear maps have a matrix representation. (Analogous examples occur for infinite dimensional 2-vector spaces).

Lemma

If \mathbf{B} is an essentially finite groupoid, the functor category $\Lambda(\mathbf{B}) = \text{Rep}(\mathbf{B}) := [\mathbf{B}, \mathbf{Vect}]$ is a KV 2-vector space.

The generators of $[\mathbf{B}, \mathbf{Vect}]$ are irreducible reps - labeled by $([b], V)$, where $[b] \in \underline{\mathbf{B}}$ and V an irreducible rep of $\text{Aut}(b)$.

Theorem

If \mathbf{X} and \mathbf{B} are essentially finite groupoids, a functor $f : \mathbf{X} \rightarrow \mathbf{B}$ gives two 2-linear maps:

$$f^* : \Lambda(\mathbf{B}) \rightarrow \Lambda(\mathbf{X})$$

with $f^*F = F \circ f$ and (the restricted representation along f)

$$f_* : \Lambda(\mathbf{X}) \rightarrow \Lambda(\mathbf{B})$$

the induced representation of F along f . Furthermore, f_* is the two-sided adjoint to f^* .

In fact, the LEFT adjoint map f_* acts by:

$$f_*(F)(b) \cong \bigoplus_{f(x) \cong b} \mathbb{C}[Aut(b)] \otimes_{\mathbb{C}[Aut(x)]} F(x)$$

This is a (left) *Kan extension* of the functor F along f .

There is also a RIGHT adjoint (right Kan extension):

$$f_!(F)(b) \cong \bigoplus_{[x] | f(x) \cong b} \text{hom}_{\mathbb{C}[Aut(x)]}(\mathbb{C}[Aut(b)], F(x))$$

There is the canonical *Nakayama isomorphism*:

$$N_{(f,F,b)} : f_!(F)(b) \rightarrow f_*(F)(b)$$

given by the *exterior trace map* (which uses a modified group average in each factor):

$$N : \bigoplus_{[x] \mid f(x) \cong b} \phi_x \mapsto \bigoplus_{[x] \mid f(x) \cong b} \frac{1}{\#Aut(x)} \sum_{g \in Aut(b)} g \otimes \phi_x(g^{-1})$$

Under this identification we get that f^* and f_* are ambidextrous adjoints.

Call the adjunctions in which f_* is left or right adjoint to f^* the *left and right adjunctions* respectively. We want to use the counit for the right adjunction, the evaluation map:

$$\eta_R(G)(x) : v \mapsto \bigoplus_{y|f(y)\cong x} (g \mapsto g(v))$$

and the unit for the left adjunction, which is determined by the action:

$$\epsilon_L(G)(x) : \bigoplus_{[y]|f(y)\cong x} g_y \otimes v \mapsto \sum_{[y]|f(y)\cong x} f(g_y)v$$

These define maps between $F(x)$ and $f_*f^*F(x)$.

(Note: there are canonical inner products around which make these maps *linear adjoints*.)

Definition

Define the 2-functor Λ as follows:

- Objects: $\Lambda(\mathbf{B}) = \text{Rep}(\mathbf{B}) := [\mathbf{B}, \mathbf{Vect}]$
- Morphisms $\Lambda(X, s, t) = t_* \circ s^* : \Lambda(\mathbf{a}) \longrightarrow \Lambda(\mathbf{B})$
- 2-Morphisms: $\Lambda(Y, \sigma, \tau) = \epsilon_{L, \tau} \circ N \circ \eta_{R, \sigma} : (t)_* \circ (s)^* \rightarrow (t')_* \circ (s')^*$

$\Lambda(X, s, t)$ can be represented by the matrix with coefficients:

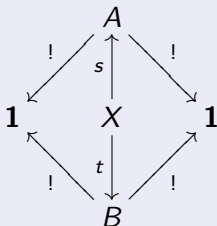
$$\begin{aligned} \Lambda(X, s, t)_{([a], V), ([b], W)} &= \text{hom}_{\text{Rep}(\text{Aut}(b))}(t_* \circ s^*(V), W) \\ &\simeq \bigoplus_{[x] \in \underline{(s, t)^{-1}([a], [b])}} \text{hom}_{\text{Rep}(\text{Aut}(x))}(s^*(V), t^*(W)) \end{aligned}$$

This is an *intertwiner space* for the groupoid representations. The 2-morphisms give (component-wise) linear maps between intertwiner spaces.

In the case where source and target are $\mathbf{1}$, there is only one basis object in $\Lambda(\mathbf{1})$ (the trivial representation), so the 2-linear maps are represented by a single vector space. Then it turns out:

Theorem

Restricting to $\text{hom}_{\text{Span}_2(\mathbf{Gpd})}(\mathbf{1}, \mathbf{1})$:



where $\mathbf{1}$ is the (terminal) groupoid with one object and one morphism, Λ on 2-morphisms is just the degroupoidification functor D of Baez and Dolan.

Theorem

For any finite group G , the 2-functor

$$Z_G = \Lambda \circ \mathcal{A}_G$$

is an extended TQFT.

That is, a cobordism becomes:

$$[\mathcal{A}_G(X), \mathbf{Vect}]^{\Lambda(\mathcal{A}_G(S), (i_X)^*, (i_Y)^*)} [\mathcal{A}_G(Y), \mathbf{Vect}]$$

and similarly for 2-morphisms.

Corollary

$Z_G = \Lambda \circ \mathcal{A}_G$ gives the Dijkgraaf-Witten model when $n = 3$, for closed manifolds.

For example, if $B = S^1$, a G -connection g is determined by one holonomy in G . Then G acts by gauge transformations, via the conjugation action, so $h : g \rightarrow g'$ is $h \in G$, such that $g' = hgh^{-1}$. So:

$$\mathcal{A}_G(S^1) \simeq G // G$$

which has:

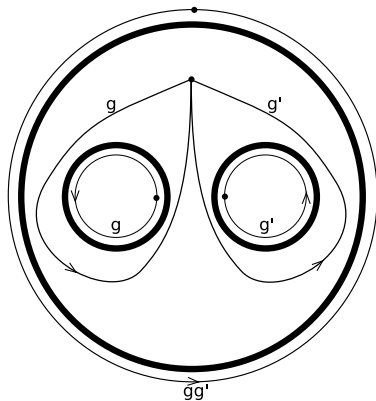
- Objects: elements of G
- Morphisms: conjugacy relations $h : g \rightarrow g'$

So we have the 2-vector space of G -equivariant functors into **Vect**, and thus

$$Z_G(S^1) = [\mathcal{A}_G(S^1), \mathbf{Vect}] \simeq \text{Rep}(G // G)$$

The irreducible (basis) objects of $Z_G(S^1)$ are then labelled by a choice of conjugacy class $[g] \in \underline{G}$ and a representation $V \in \text{Rep}(\text{Stab}(g))$.

Suppose $S : S^1 + S^1 \rightarrow S^1$ is the “pair of pants”, showing two “particles” fusing into one.



Then we have the diagram:

$$\begin{array}{ccc} & (G \times G) // G & \\ \Delta \swarrow & & \searrow m \\ (G // G)^2 & & G // G \end{array} \quad (1)$$

Where the map Δ leaves connections fixed, and acts as the diagonal on gauge transformations; and m is the multiplication map.

Note: a state over $([g], [g'])$ will be transported to one with nontrivial representations over all $[gg']$ for any representatives of $[g], [g']$.

Generalization 1: Z_G for G a compact Lie group (uses measured groupoids)

For quantum mechanics, the classical configuration space S is usually not discrete.

Minimally, (S, μ) a measure space, and $L^2(S, \mu)$ is the Hilbert space for the corresponding quantum system. Interesting cases occur when S is a manifold and μ comes from a volume form.

Duplicating the above requires some changes:

- Ambi-adjunction requires **Hilb** instead of **Vect** in infinite-dim setting
- Direct sums become direct integrals - which are not (co)limits
- Push-forward is not just Kan extension of functors

Any 2-linear map $T : \mathbf{Vect}^n \rightarrow \mathbf{Vect}^m$ is naturally isomorphic to a map acting by an $m \times n$ matrix:

$$\begin{pmatrix} V_{1,1} & \cdots & V_{1,n} \\ \vdots & & \vdots \\ V_{m,1} & \cdots & V_{m,n} \end{pmatrix} \begin{pmatrix} W_1 \\ \vdots \\ W_n \end{pmatrix} = \begin{pmatrix} \bigoplus_{i=1}^n V_{1,i} \otimes W_i \\ \vdots \\ \bigoplus_{i=1}^n V_{m,i} \otimes W_i \end{pmatrix}$$

When the entries are finite-dimensional vector spaces, this explains why T has a two-sided adjoint.

T^* is the $n \times m$ matrix with $(T^*)_{i,j} = (T_{i,j})^*$, the dual of the corresponding entry in the transpose of T . The adjoint is 2-sided because $(V_{i,j})^{**} \cong V_{i,j}$ is canonical: the category **FinVect** is *reflexive*.

This isn't true for infinite-dimensional vector spaces, but it is for Hilbert spaces (**Hilb** is reflexive). So to stay closed under composition, in infinite-dimensions, 2-Vector spaces must be generalized to 2-Hilbert spaces.

Consider categories like:

Definition

If (X, μ) is a measurable space $\mathbf{Meas}(X)$ is the category with:

- Objects: *measurable fields of Hilbert spaces* on (X, \mathcal{M}) : i.e. X -indexed families of Hilbert spaces \mathcal{H}_x with a Hilbert space of *measurable sections* (satisfying certain properties)
- Morphisms: *measurable fields of bounded linear maps* between Hilbert spaces, $f_x : \mathcal{H}_x \rightarrow \mathcal{K}_x$ so that $\|f\|$ (the operator norm of f) is measurable.

This is the equivalent of a *measurable function*. Imposing that sections and norms be L^2 condition gives a categorification of $L^2(X, \mu)$.

Definition

There is a locale MX whose “open sets” are the measurable sets of X , and whose morphisms are inclusions *up to almost everywhere*. This becomes a Grothendieck site where an “open cover” is a usual cover, *up to almost everywhere*.

Then a *measurable sheaf of Hilbert spaces* is a sheaf of Hilbert spaces on MX , and these form a category $MSh(X)$.

Theorem (Wendt)

The category of measurable sheaves $MSh(X)$ is equivalent to the internal category $Hilb[Sh(X)]$ of Hilbert spaces in the topos $Sh(X)$ of (set-valued) sheaves on MX .

Theorem

A measurable field of Hilbert spaces on (X, μ) determines a measurable sheaf by direct integration: given a measurable $U \subset X$, this assigns

$$\int_U^{\oplus} d\mu(x) \mathcal{H}_x$$

where the direct integral is a Hilbert space of sections with inner product

$$\langle \phi, \psi \rangle = \int_U d\mu(x) \langle \phi_x, \psi_x \rangle$$

This is the equivalent of the matrix of vector spaces for a 2-linear map. It is still a conjecture that all suitable functors are of this form.

Question: How do such functors arise?

Definition

A *disintegration* between two measure spaces consists of:

- A measurable function $f : (X, \mathcal{M}, \mu) \rightarrow (Y, \mathcal{N}, \nu)$
- A family $(X_y, \mathcal{M}_y, \mu_y)_{y \in Y}$ where:
 - ▶ $X_y = f^{-1}(y)$
 - ▶ $\mathcal{M}_y = \{A \cap X_y \mid A \in \mathcal{M}\}$
 - ▶ μ_y is a measure on X_y

satisfying some obvious properties.

Theorem (Wendt)

Given a disintegration $f : (X, \mu) \rightarrow (B, \nu)$, there is an adjoint pair of functors

$$MSh(X) \begin{array}{c} \xrightarrow{f^*} \\ \xleftarrow{f_*} \end{array} MSh(Y)$$

We need (groupoid-)equivariant version of this theorem.

Definition

A measurable groupoid is a groupoid internal to **Msble**, the category of measurable spaces and measurable functions.

Definition

If $\mathcal{G} = (G_0, G_1)$ is a measurable groupoid, a **groupoid measure** on \mathcal{G} consists of:

- A measure μ on the space of objects
- A (measurable, left) *Haar system*: for each $x \in G_0$, a measure ν_x on the space of morphisms into x , $t^{-1}(x)$ such that
 - ▶ the choice ν_x is measurable: for any measurable function $f : G_1 \rightarrow \mathbb{C}$, the function

$$x \mapsto \int_{t^{-1}(x)} f(g) d\nu_x(g) \quad (2)$$

is measurable

- ▶ the ν_x are left-invariant: for any $g \in G_1$, and measurable $f : G_1 \rightarrow \mathbb{C}$

$$\int f(gh) d\nu_{s(g)}(h) = \int f(h) d\nu_{t(g)}(h) \quad (3)$$

To define Λ for measure groupoids, we again want:

$$\Lambda(G) = \text{Rep}(G)$$

A representation ρ of a measure groupoid $s, t : G_1 \rightarrow G_0$ is defined on a measurable field of Hilbert spaces \mathcal{H} on G_0 . It gives a functor $R : \mathbf{G} \rightarrow \mathbf{Hilb}$ with $R(x) = \mathcal{H}_x$, the fibre at each $x \in M$, and an isomorphism $R(g)$ for each $g : x \rightarrow y$. (But not all functors are *measurable* representations).

Definition

$\text{Rep}(\mathbf{G})$, the **category of representations** of \mathbf{G} , has

- *Objects*: Measurable representations of \mathbf{G}
- *Morphisms*: Intertwiners: i.e. measurable natural transformations between functors $n : \rho \rightarrow \rho'$

(A natural transformation is measurable when it determines a measurable field of linear maps over G_0 .)

Theorem

A representation of G on a measurable field \mathcal{H} of Hilbert spaces determines an equivariant sheaf of Hilbert spaces by direct integration.

Then we hope to have the following:

Proposition

The category $\text{Rep}(G)$ is equivalent to the internal category $\text{Hilb}[\text{EMSh}(G)]$ of Hilbert spaces in the topos of equivariant measurable sheaves on MG_0 .

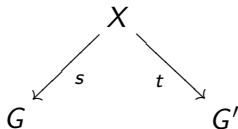
And

Proposition

Given a disintegrating functor $f : G \rightarrow G'$ between measure groupoids, there is a (bi-)adjoint pair of functors

$$\text{EMSh}(G) \begin{array}{c} \xrightarrow{f^*} \\ \xleftrightarrow{\quad} \\ \xleftarrow{f_*} \end{array} \text{EMSh}(G')$$

Given the above, we would define Λ as before, so that a span



has

$$\Lambda(X, s, t) = t_* \circ s^* : \Lambda(G) \longrightarrow \Lambda(G')$$

and for a 2-cell $Y : X \rightarrow X'$ given by

$$\Lambda(Y, \sigma, \tau) = \epsilon_{L, \tau} \circ N \circ \eta_{R, \sigma} : (t)_* \circ (s)^* \rightarrow (t')_* \circ (s')^*$$

using the analog of the Nakayama isomorphism:

$$N : \int_{[x] | f(x) \cong b}^{\oplus} \phi_x \mapsto \int_{[x] | f(x) \cong b}^{\oplus} \frac{1}{\text{vol}(\text{Aut}(x))} \int_{g \in \text{Aut}(b)} g \otimes \phi_x(g^{-1})$$

Applying the above to ETQFT: follow the same prescription:

Example

Interesting case is $G = SU(2)$. The topology generates measurable sets to make $SU(2)$ a regular Borel space, with Haar measure μ .

The (measurable) groupoid

$$\mathcal{G} = Z_{SU(2)}(S^1) = SU(2) // SU(2)$$

gets a groupoid measure:

- **Measure:** $Ob(\mathcal{G}) = SU(2)$ gets the Haar measure μ
- **Haar System:** For $g \in Ob(\mathcal{G})$, we always have $t^{-1}(g) \cong SU(2)$, which also gets $\nu_g = \mu$

(Note: $vol(G // G) = 1$, as we've fixed normalization of μ)

We can get reps of \mathcal{G} by integrating those indexed by $([g], V)$ for $g \in SU(2)$ and V an irrep of $Stab(g)$ ($SU(2)$ or $U(1)$).

Cobordisms of 2 or 3 dimensions are trickier:

For connected cobordisms, all groupoids in our construction are equivalent to ones of the form $\mathcal{A}_G(X) // G$.

So we can always take $\nu_x = \mu$, Haar measure on G . But:

- There is a canonical measure on $\mathcal{A}_G(B)$ for 2-manifold S , the *Goldman measure*... but this is nontrivial!
- There is no *canonical* measure on $\mathcal{A}_G(M)$ for 3-manifold M !

To assign measures to $\mathcal{A}_G(X)$ in dimension 3 or higher, we must *change the cobordism category*.

Need cobordisms to be decorated with extra data, sufficient to determine a measure. (e.g. specified paths which determine a presentation of the fundamental groupoid)

The construction for Λ can also be described using:

- $Rep(\mathbf{B}) \mapsto$ Category of reps of von Neumann algebra associated to \mathbf{B} (including groupoid algebras)
- 2-linear maps represented by *Hilbert bimodules*, given by induction and restriction
- Natural transformations represented by bimodule maps

This relates to a conjecture (Baez, Baratin, Freidel, Wise) that *infinite-dimensional 2-Hilbert spaces* are representation categories for v.N.-algebras.

“Physically”, this describes the quantum mechanics of systems with boundary.

Generalization 2: Higher gauge theory - for a 2-group \mathcal{G} , define a 3-functor $Z_{\mathcal{G}} : \mathbf{nCob}_3 \rightarrow \mathbf{3Vect}$.

Sketch:

Definition

Fixing a 2-group \mathcal{G} , the contravariant 2-functor

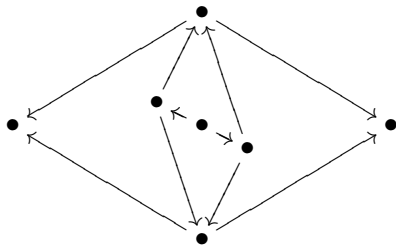
$$\mathcal{A}_{\mathcal{G}}^{(2)} = 2\text{Fun}[\Pi_2(-), \mathcal{G}]$$

assigns to a manifold M the 2-groupoid $\mathcal{A}_{\mathcal{G}}^{(2)}(M)$ with:

- Objects: 2-functors (“2-connections”)
- Morphisms: natural transformations (“gauge transformations”)
- 2-Morphisms: modifications (...)

(and, to smooth functions, the induced maps).

There's an induced map $Span_3(\mathbf{ManCorn}) \rightarrow Span_3(\mathbf{2Gpd})$, where $Span_3(-)$ has, as 3-morphisms, equivalence classes of diagrams shaped like:



(4)

Composition is again by weak pullback. (Note that 2-morphisms and 3-morphisms of $\mathbf{2Gpd}$ can appear in $Span_3(\mathbf{2Gpd})$ by weakening the assumption that this commutes.)

As before, $nCob_3$ lives in $Span^3(\mathbf{ManCorn})$ (cubical, but can be intimidated into being globular if desired).

We would like to define a 3-functor

$$\Lambda^{(2)} : \mathit{Span}_3(\mathbf{2Gpd}) \rightarrow \mathbf{3Vect}$$

Then assuming $\Lambda^{(2)}$ is well-defined, we should obtain an extended TQFT 3-functor:

$$Z_{\mathcal{G}} = \Lambda^{(2)} \circ \mathcal{A}_{\mathcal{G}}^{(2)} : \mathbf{nCob}_3 \rightarrow \mathbf{3Vect}$$

For $\mathcal{X} \in \mathbf{2Gpd}$, we expect to get:

$$\Lambda^{(2)}(\mathcal{X}) = \mathit{Rep}(\mathcal{X}) = \mathit{2Fun}(\mathcal{X}, \mathbf{2Vect})$$

There should be, for each $\mathcal{F} : \mathcal{X} \rightarrow \mathcal{Y}$, an adjoint pair

$$\mathcal{F}^* : \text{Rep}(\mathcal{Y}) \rightarrow \text{Rep}(\mathcal{X})$$

and

$$\mathcal{F}_* : \text{Rep}(\mathcal{X}) \rightarrow \text{Rep}(\mathcal{Y})$$

where the induced representation functor \mathcal{F}_* is given by 2-Kan extension along \mathcal{F} .

To prove: It should be biadjoint. Moreover, to get 3-morphisms, the unit and counit ϵ_L, η_R should themselves be biadjoint!

Eventually: One hopes this pattern will repeat with representations of n -groupoids for all n . Must deal with slight trickiness of **nVect**.

