

2-Linearization In Physics and Topology

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Motivation: Categorify a quantum mechanical description of states and processes.

Applications Foundational physics such as quantum harmonic oscillator; Witten-type ETQFT (help interpret physical examples).

Categorification involves replacing set-based structures with category-based structures. That is, by replacing the category **Set** with the 2-category **Cat** (or **SmallCat**). There are two obvious approaches to how the original structure reappears (apart from “by analogy”):

- 1 “Quotient”: from the object/morphism level (Grothendieck ring - e.g. categorified \mathfrak{sl}_2)
- 2 “Substructure”: At the morphism/2-morphism level (the “microcosm principle”)

There are more possibilities when going to n -categories.

The following is an example of the last type:

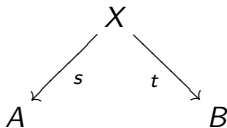
Theorem

There is a 2-functor (“2-linearization”):

$$\Lambda : \text{Span}_2(\mathbf{Gpd}) \rightarrow \mathbf{2Vect}$$

This is a categorification (sense 2) of the “degroupoidification” functor $D : \text{Span}_1(\mathbf{Gpd}) \rightarrow \mathbf{Vect}$ of Baez and Dolan (which itself gives an example of sense 1!)

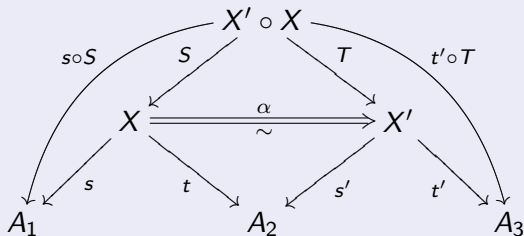
A *span* in a (n -)category \mathbf{C} is a diagram:



The bicategory $\text{Span}_2(\mathbf{Gpd})$ (similar for any 2-category with weak pullbacks) has:

Definition

- **Objects:** Groupoids
- **Morphisms:** Spans of groupoids
- Composition defined by *weak pullback*:



- **2-Morphisms** : isomorphism classes of *spans of span maps*
- monoidal structure from the product in \mathbf{Gpd} , monoidal unit 1

(The category $\text{Span}(\mathbf{Gpd})$ takes spans up to span-isomorphism)

Definition

$$D : \text{Span}(\mathbf{Gpd}) \rightarrow \mathbf{Vect}$$

with $D(G) = \mathbb{C}(\underline{G})$,

$$D(X)(f)([b]) = \sum_{[x] \in \underline{t^{-1}(b)}} \frac{\# \text{Aut}(b)}{\# \text{Aut}(x)} [f(s(x))]$$

This amounts to multiplication by a matrix $D(X)$ with

$$D(X)_{([a],[b])} = |(s, t)^{-1}(a, b)|$$

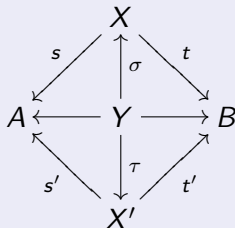
using *groupoid cardinality* (which can be interpreted as an inner product in a canonical way).

(*Note:* Compare gpd. cardinality to role of Euler characteristic in geom. representation theory.)

D is a “quotient-style” decategorification map.

Definition (Part 2)

The **2-morphisms** of $\text{Span}_2(\mathbf{Gpd})$ are (iso. classes of) spans of *span maps*:



Composition is by weak pullback taken up to isomorphism.

(Often one just uses span maps: here, we want 2-morphisms as well as morphisms to have *adjoints*, so much use these.)

Definition

2Vect is the 2-category of *Kapranov-Voevodsky* 2-vector spaces, which consists of:

- Objects: **Kapranov–Voevodsky 2-vector spaces**: \mathbb{C} -linear finitely semisimple additive category (one generated by simple objects x , where $\text{hom}(x, x) \cong \mathbb{C}$).
- Morphisms: **2-linear maps**: \mathbb{C} -linear (hence additive) functor.
- 2-Morphisms: Natural transformations between 2-linear maps

Note: **2Vect** is a monoidal 2-category with the Deligne product and unit **Vect**.

Theorem (KV)

*Any KV 2-vector space is equivalent to **Vect**^k for some k. Any 2-linear map is then naturally isomorphic to one given by a matrix of vector spaces (and matrix multiplication using \otimes and \oplus). Any natural transformation of 2-linear maps is then given by a matrix of componentwise linear maps.*

Lemma

If \mathbf{B} is an essentially finite groupoid, the functor category $\Lambda(\mathbf{B}) = [\mathbf{B}, \mathbf{Vect}]$ is a KV 2-vector space.

Note: If the automorphism groups of (isomorphism classes of) objects of \mathbf{B} are B_1, \dots, B_n , then we have

$$[\mathbf{B}, \mathbf{Vect}] \cong \prod_j \text{Rep}(B_j)$$

So the “basis elements” (simple objects) in $[\mathbf{B}, \mathbf{Vect}]$ are labeled by $([b], V)$, where $[b] \in \mathbf{B}$ and V an irreducible rep of $\text{Aut}(b)$.

Theorem

If \mathbf{X} and \mathbf{B} are essentially finite groupoids, a functor $f : \mathbf{X} \rightarrow \mathbf{B}$ gives two 2-linear maps:

$$f^* : \Lambda(\mathbf{B}) \rightarrow \Lambda(\mathbf{X})$$

namely composition with f , with $f^*F = F \circ f$ and

$$f_* : \Lambda(\mathbf{X}) \rightarrow \Lambda(\mathbf{B})$$

called “pushforward along f ”. Furthermore, f_* is the two-sided adjoint to f^* (i.e. both left-adjoint and right-adjoint).

In fact, the adjoint map f_* acts by:

$$f_*(F)(b) \cong \bigoplus_{f(x) \cong b} \mathbb{C}[Aut(b)] \otimes_{\mathbb{C}[Aut(x)]} F(x)$$

This is the left adjoint. But there is also a right adjoint:

$$f_!(F)(b) \cong \bigoplus_{[x] | f(x) \cong b} \text{hom}_{\mathbb{C}[Aut(x)]}(\mathbb{C}[Aut(b)], F(x))$$

The *Nakayama isomorphism* is a canonical isomorphism between these (in particular: it defines an isomorphism even over base rings other than \mathbb{C}). It gives maps:

$$N_{(f,F,b)} : f_!(F)(b) \rightarrow f_*(F)(b)$$

given by the *exterior trace map* in each factor of the sum (which uses a modified group average):

$$N : \bigoplus_{[x] \mid f(x) \cong b} \phi_x \mapsto \bigoplus_{[x] \mid f(x) \cong b} \frac{1}{\#Aut(x)} \sum_{g \in Aut(b)} g \otimes \phi_x(g^{-1})$$

Under this identification, the left and right adjoints are isomorphic. By composing units/counters with N , we get that f^* and f_* are ambidextrous adjoints.

(Note: In general, $Span_2(\mathbf{C})$ will be the universal 2-category for which morphisms in \mathbf{C} have ambidextrous adjoints. We want to preserve this property.)

Call the adjunctions in which f_* is left or right adjoint to f^* the *left and right adjunctions* respectively. We want to use the counit for the left adjunction, which is the evaluation map:

$$\begin{aligned} \eta_R(G)(x) : G(x) &\rightarrow \bigoplus_{y|f(y)\cong x} \text{hom}_{\mathbb{C}[Aut(x)]}(\mathbb{C}[Aut(y)], G(x)) \\ v &\mapsto \bigoplus_{y|f(y)\cong x} (g \mapsto g(v)) \end{aligned}$$

and the unit for the right adjunction, which just uses the action:

$$\begin{aligned} \epsilon_L(G)(x) : \bigoplus_{[y]|f(y)\cong x} \mathbb{C}[Aut(x)] \otimes_{\mathbb{C}[Aut(y)]} f^* G(x) &\rightarrow G(x) \\ \bigoplus_{[y]|f(y)\cong x} g_y \otimes v &\mapsto \sum_{[y]|f(y)\cong x} f(g_y)v \end{aligned}$$

Definition

Define the 2-functor Λ as follows:

- Objects: $\Lambda(\mathbf{B}) = \text{Rep}(\mathbf{B}) := [\mathbf{B}, \mathbf{Vect}]$
- Morphisms $\Lambda(X, s, t) = t_* \circ s^* : \Lambda(\mathbf{a}) \longrightarrow \Lambda(\mathbf{B})$
- 2-Morphisms: $\Lambda(Y, \sigma, \tau) = \epsilon_{L, \tau} \circ N \circ \eta_{R, \sigma} : (t')_* \circ (s')^* \rightarrow (t)_* \circ (s)^*$

Picking basis elements $([a], V) \in \Lambda(A)$, and $([b], W) \in \Lambda(B)$, we get that $\Lambda(X, s, t)$ is represented by the matrix with coefficients:

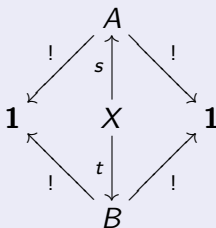
$$\begin{aligned} \Lambda(X, s, t)_{([a], V), ([b], W)} &= \text{hom}_{\text{Rep}(\text{Aut}(b))}(t_* \circ s^*(V), W) \\ &\simeq \bigoplus_{[x] \in \underline{(s, t)^{-1}([a], [b])}} \text{hom}_{\text{Rep}(\text{Aut}(x))}(s^*(V), t^*(W)) \end{aligned}$$

This is an intertwiner space, given by the *analog* of an inner product. The 2-morphisms give linear maps between intertwiner spaces, which can also be interpreted as a “pull-push” operation.

In the case where source and target are $\mathbf{1}$, there is only one basis object in $\Lambda(\mathbf{1})$ (the trivial representation), so the 2-linear maps are represented by a single vector space. Then it turns out:

Theorem

Restricting to $\text{hom}_{\text{Span}_2(\mathbf{Gpd})}(\mathbf{1}, \mathbf{1})$:



where $\mathbf{1}$ is the (terminal) groupoid with one object and one morphism, Λ on 2-morphisms is just the degroupoidification functor D .

The groupoid cardinality comes from the modified group average in N .

2-Linearized Physics

“Physical” applications arise because groupoids provide a good way of thinking about local symmetry, generalizing the *action groupoid* $S//G$ associated to a G -action on S .

In a span $A \leftarrow X \rightarrow B$, the groupoid X will represent a space of *histories*; s and t pick the starting and terminating *configuration* in spaces A and B .

This setup is how we “do physics in” the monoidal (2-)category $\mathbf{Span}_2(\mathbf{Gpd})$. The functors D and Λ will give a description of physics in \mathbf{Vect} (really, \mathbf{Hilb} since there is a canonical inner product), and $\mathbf{2Vect}$ respectively (ditto).

The span $\mathbf{Vect} \leftarrow \mathbf{Span}_2(\mathbf{Gpd}) \rightarrow \mathbf{2Vect}$ provides a way to “categorify quantum mechanics”.

Example

An Extended TQFT (ETQFT) is a (weak) monoidal 2-functor

$$Z : \mathbf{nCob}_2 \rightarrow \mathbf{2Vect}$$

where \mathbf{nCob}_2 has

- **Objects:** $(n - 2)$ -dimensional manifolds
- **Morphisms:** $(n - 1)$ -dimensional cobordisms (manifolds with boundary, with ∂M a union of source and target objects)
- **2-Morphisms:** n -dimensional cobordisms with corners

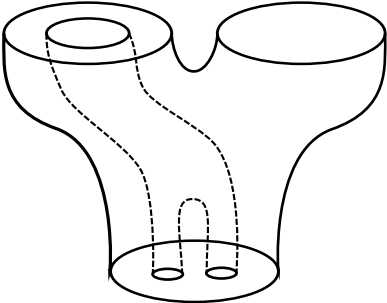
One construction uses *gauge theory*, for gauge group G (here a finite group). Given M , the groupoid $\mathcal{A}_0(M, G) = \text{hom}(\pi_1(M), G) // G$ has:

- **Objects:** Flat connections on M
- **Morphisms** Gauge transformations

Then $\mathcal{A}_0(-, G) : \mathbf{nCob}_2 \rightarrow \text{Span}_2(\mathbf{Gpd})$, and there is an ETQFT $Z_G = \Lambda \circ \mathcal{A}_0(-, G)$.

It happens to give the *Dijkgraaf-Witten model* when $n = 3$.

This relies on the fact that cobordisms in \mathbf{nCob}_2 can be transformed into products of cospans:

\mathbf{nCob}_2	$Span^2(Top)$
	$ \begin{array}{ccccc} S^1 & \xrightarrow{i_A} & (A \amalg D) & \xleftarrow{i'_A \otimes i_D} & S^1 \amalg S^1 \\ \downarrow i_1 & & \downarrow \iota_1 & & \downarrow i_2 \\ Y & \xrightarrow{\iota_3} & M & \xleftarrow{\iota_4} & Y \\ \uparrow i_2 & & \uparrow \iota_2 & & \uparrow i_1 \\ S^1 \amalg S^1 & \xrightarrow{i_2} & Y & \xleftarrow{i_1} & S^1 \end{array} $

Then $\mathcal{A}_0(-, G)$ maps these into $Span^2(\mathbf{Gpd})$.

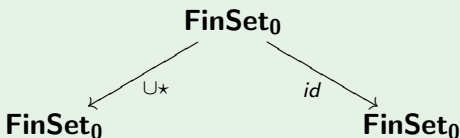
Example

In the case where $\mathbf{A} = \mathbf{B} = \mathbf{FinSet}_0$ (equivalently, the symmetric groupoid $\coprod_{n \geq 0} \Sigma_n$ - note no longer finite), we find

$$D(\mathbf{FinSet}_0) = \mathbb{C}[[t]]$$

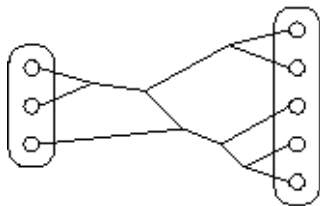
where t^n marks the basis element for object $[n]$. This gets a canonical inner product and can be treated as the Hilbert space for the *quantum harmonic oscillator* (“Fock Space”).

The operators $\mathbf{a} = \partial_t$ and $\mathbf{a}^\dagger = M_t$, generate the *Weyl algebra* of operators for the QHO. These are given under D by the span A :



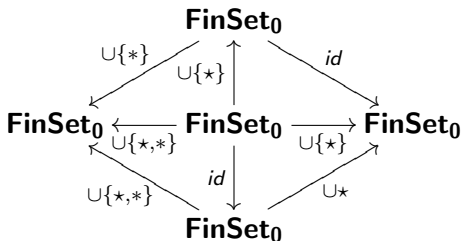
and its dual A^\dagger . Composites of these give a categorification of operators explicitly in terms of *Feynman diagrams*.

Such composites are described in terms of groupoids whose objects look like this:



The source and target maps for the span pick the set of start and end points. The morphisms of the groupoid are graph symmetries. Degroupoidification D calculates operators which (after small modification involving $U(1)$ -labels) agree with the usual Feynman rules for calculating amplitudes.

An ongoing project (with Jamie Vicary) is to study the 2-categorical version of this picture. There are analogs of creation and annihilation operators in other *hom*-categories than $hom(1, 1)$:



This is a 2-morphism $\alpha_A : A \rightarrow AAA^\dagger$ creates a “creation/annihilation pair” at the 1-morphism level.

Composites of these act as *rewrite rules* on the Feynman diagrams like those seen previously (now with “boundary” maps).

Toward Real QFT

Both the QHO and TQFT are “baby” models of real QFT, which is much harder.

One ingredient: the construction for Λ can be extended to measure-groupoids (e.g. derived from compact Lie groups w/ Haar measure), using:

- **Vect** \mapsto **Hilb** (ambiadjoint uses double-dual isomorphism)
- $\text{Rep}(\mathbf{B})$ \mapsto Category of reps of von Neumann algebra associated to \mathbf{B}
- 2-linear maps represented by *Hilbert bimodules*
- Direct sum \mapsto direct integral
- Groupoid cardinality \mapsto volume of groupoid (c.f. Weinstein)

This relates to a conjecture of Baez et. al. that *infinite-dimensional 2-Hilbert spaces* are equivalent to representation categories for v.N.-algebras.

